

Mathematical Tools for Parameters Identification

Amin Boumenir and Vu Kim Tuan
Department of Mathematics
University of West Georgia
boumenir@westga.edu

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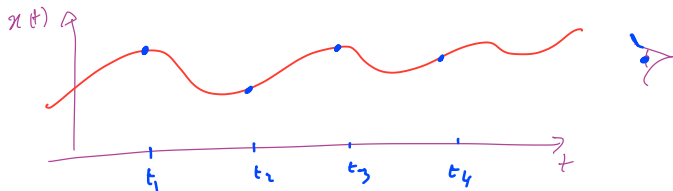
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- Hauer, J.F.; Demeure, C.J.; Scharf, L.L. (1990). "Initial results in Prony analysis of power system response signals". IEEE Transactions on Power Systems. 5: 80-89.
- Carriere, R.; Moses, R.L. (1992). "High resolution radar target modeling using a modified Prony estimator". IEEE Transactions on Antennas and Propagation. 40: 13-18.

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- $n = 1.66 + 0.25l \Rightarrow$ Leukemia,

Detecting Leukemia: Using Computational Electromagnetics

David Colton and Peter Monk

IEEE Computational Science Engineering archive

Volume 2 Issue 4, December 1995 Pages 46-52

IEEE Computer Society Press Los Alamos, CA, USA

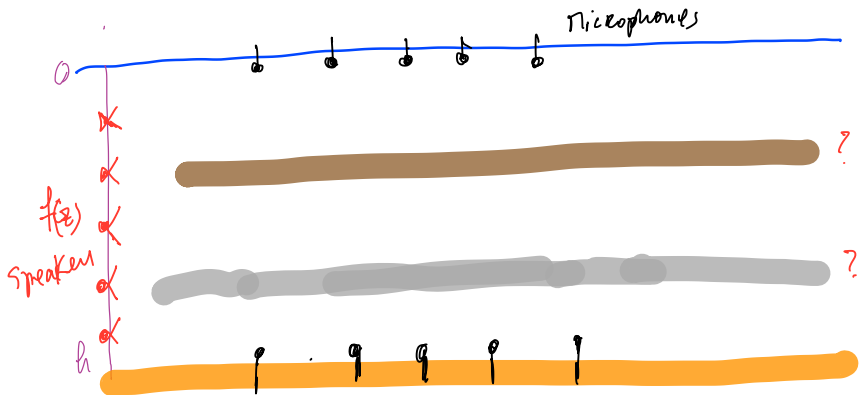
Habib Ammari

An Introduction to Mathematics of Emerging Biomedical Imaging

Springer Science 2008

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- Fadhel Al-Musallam, A.B.,
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SIAM J. Appl. Math. 71, (2011), 972-982.

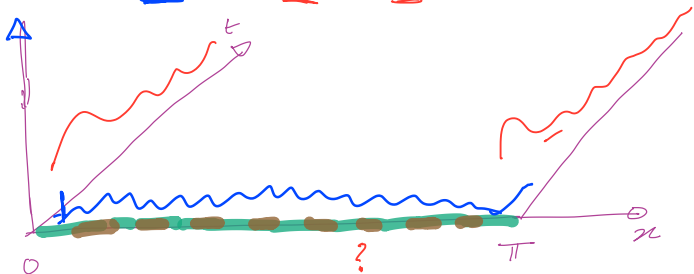
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- Recovery of the heat coefficient by two measurements, Inverse Problems and Imaging, 5, No. 4, (2011), 775-791

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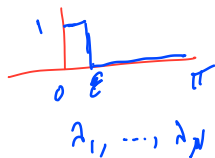
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- $c_n(f) = \int_0^\pi f(x) \varphi(x, \lambda_n) dx \neq 0$ how to choose f ?

$$f(x) = x^\alpha$$

$$-\frac{1}{2} < \alpha < 0$$

Large $\lambda_n \rightarrow \infty$



Laplace transform to identify

- How to extract spectral data :: $g(t) = \sum_{n \geq 1} a_n e^{-\lambda_n t}$
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- then consider $g(t) - a_1 e^{-\lambda_1 t} = \sum_{n \geq 2} a_n e^{-\lambda_n t}$
- Repeat the procedure to get a_2, λ_2, \dots

- $\partial_t u = \Delta u + qu$

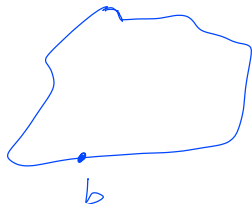
- $\partial_t u = \Delta u + qu$
- Reconstruct $q \in C(\Omega)$

$$[D] \begin{cases} u_t(x, t) = \Delta u(x, t) + q(x)u(x, t) & x \in \Omega \subseteq \mathbb{R}^d, \quad d \geq 2 \\ u(x, t) = 0 \quad \text{for } x \in \partial\Omega \\ u(x, 0) = f(x) \in L^2(\Omega) \end{cases}$$

from the map

$$f \longrightarrow \{\partial_\nu u(b, k)\}_{k \geq 1}$$

where $b \in \partial\Omega$



Borg type results on $\Delta + q$

$$\begin{cases} \Delta \varphi_n(x) + q(x)\varphi_n(x) = -\lambda_n \varphi_n(x) \\ \varphi_n(x) = 0 \text{ on } \partial\Omega \end{cases}$$

Result: Nachman Uhlmann, Sylvester,

$\left\{ \lambda_n, \frac{\partial}{\partial \nu} \varphi_n(x) \mid x \in \partial\Omega \right\}_{n \geq 1}$ determines q uniquely.

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- A.B. and Vu Kim Tuan, Inverse problems for multidimensional heat equations by measurements at a single point on the boundary, *Numer. Funct. Anal. Optim.* 30(2009), no. 11-12, 1215-1230

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① let $\varphi_n(x)$ be the eigenfunctions

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$$\begin{cases} \Delta \phi_n(x) + q(x)\phi_n(x) = -\lambda_n \phi_n(x) \\ \partial_\nu \phi_n(x) = 0 \quad \text{for } x \in \partial\Omega \end{cases} \quad \text{where } \lambda_n \rightarrow \infty.$$

- 4 $\phi_1(b) \neq 0$? not true
- 5 for any $b \in \partial\Omega$, there is $j \geq 1$ such that $\phi_j(b) \neq 0$. Thus
- 6 From $u(b, t) = \sum_n c_n(f) e^{-\lambda_n t} \phi_n(b)$ we can extract $\{\lambda_j, \phi_j(b)c_j(\psi_k)\}_{k \geq 1} \Rightarrow \phi_j(x)$ in $C(\Omega)$
- 7 $\Delta \phi_j(x) + q(x)\phi_j(x) = -\lambda_j \phi_j(x) \Rightarrow$ reconstruct $q(x)$ a.e.

n coefficients in the the Self-Adjoint case

$$\begin{cases} u_t(x, t) = \sum_{i,j=1}^n \frac{\partial}{\partial x_i} \left(a_{ij}(x) \frac{\partial}{\partial x_j} u(x, t) \right) - q(x)u(x, t), & x \in \Omega \subset \mathbb{R}^d, \\ u(x, t) = 0, & x \in \partial\Omega, \\ u(x, 0) = \psi(x) \in L^2(\Omega), \end{cases}$$

$$-q(x)\varphi_m + \sum_{j=1}^n \frac{\partial}{\partial x_j} \left(\sum_{i=1}^n a_{ij} \right) \frac{\partial \varphi_m}{\partial x_j} + \sum_{i,j=1}^n a_{ij}(x) \frac{\partial^2 \varphi_m}{\partial x_i \partial x_j} = -\lambda_m \varphi_m.$$

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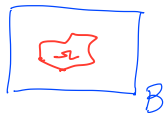
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Algebra

① For any $b \in \Omega$ there are infinitely many $\varphi_m(b) \neq 0$.

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① vanish at the origin, $\varphi_{ml}(0, \theta) = 0$ for all $m \geq 1$ for $l \in \mathbb{Z}$.

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 - 1 vanish at the origin, $\varphi_{ml}(0, \theta) = 0$ for all $m \geq 1$ for $l \in \mathbb{Z}$.
 - 2 but $\{\varphi_{0l}(0, \theta)\}_{l \in \mathbb{Z}}$ is nonzero as predicted.

Wronskians

- Reconstruct the eigenfunctions, the N -vectors $\Phi_m(x)$, with

$$\Phi_m(x) := \left(\varphi_m(x), \frac{\partial \varphi_m}{\partial x_1}(x), \frac{\partial \varphi_m}{\partial x_n}(x), \frac{\partial^2 \varphi_m}{\partial x_1^2}(x), \dots, \frac{\partial^2 \varphi_m}{\partial x_i \partial x_j}(x), \dots, \frac{\partial^2 \varphi_m}{\partial x_n^2}(x) \right) \quad (1)$$

proposition

Given a set of indices $k_1 < k_2 < \dots < k_N$, such that λ_{k_j} , $j = 1, \dots, N$, are distinct and nonzero. Then the set

$$E := \left\{ x \in \Omega : \det \left[\Phi_{k_1}^T(x), \Phi_{k_2}^T(x), \dots, \Phi_{k_N}^T(x) \right] \neq 0 \right\} \quad (2)$$

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- The main idea here is to use the constant rank theorem

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proposition

We can reconstruct the shape of Ω , the coefficients $a_{i,j}$ and q from measurements $\psi_l \rightarrow \{u(b, k)\}_{k,l \in \mathbb{N}}$ at any single point $b \in \Omega$.

A.B. ; Tuan, Vu Kim; Hoang, Nguyen

The recovery of a parabolic equation from measurements at a single point.
Evol. Equ. Control Theory 7 (2018), no. 2, 197-216.

Divergence type

- Consider the parabolic equation

$$\begin{cases} \partial_t u(x, t) = \nabla \cdot (p(x) \nabla u(x, t)), & t > 0, \quad x \in \Omega, \\ u(x, t) = 0, & x \in \partial\Omega, \\ u(x, 0) = a(x), \end{cases} \quad (3)$$

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- from a sequence of observations of the solution $\{u(b, k)\}_{k \in \mathbb{N}}$ taken at an arbitrary single point $b \in \Omega$.

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- $\text{supp} \varphi_1(x) = \overline{\Omega}$

Richter's trick,

G.R. Richter, An inverse problem for the steady state diffusion equation, *SIAM J. Appl. Math.* 41(1981), 210-221.

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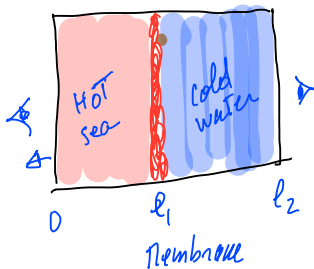
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- A.B and Tuan, Vu Kim One point recovery of a parabolic equation. *J. Math. Anal. Appl.* 463 (2018), no. 1, 161-166.

Distillation membrane

$$\left\{ \begin{array}{ll}
 \partial_t f(t, x, y) - \alpha_f \Delta f(t, x, y) + \beta_f \partial_y f(t, x, y) = 0 & t > 0, (x, y) \in \Omega_1, \\
 \partial_t p(t, x, y) - \alpha_p \Delta p(t, x, y) + \beta_p \partial_y p(t, x, y) = 0 & t > 0, (x, y) \in \Omega_2, \\
 \partial_x f(t, l_1, y) = \gamma_f (f(t, l_1, y) - p(t, l_1, y)) & t > 0, y \in (0, l_2), \\
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 \partial_x f(t, 0, y) = 0 & t > 0, y \in (0, l_2), \\
 \partial_x p(t, \delta_m + 2l_1, y) = 0 & t > 0, y \in (0, l_2), \\
 \rho(t, x, 0) = T_p(t, x) & t > 0, x \in (\delta_m + l_1, \delta_m + 2l_1), \\
 \partial_y \rho(t, x, l_2) = 0 & t > 0, x \in (\delta_m + l_1, \delta_m + 2l_1), \\
 f(t, x, 0) = T_f(t, x) & t > 0, x \in (0, l_1), \\
 \partial_y f(t, x, l_2) = 0 & t > 0, x \in (0, l_1), \\
 f(0, x, y) = f_0(x, y) & (x, y) \in \Omega_1, \\
 \rho(0, x, y) = \rho_0(x, y) & (x, y) \in \Omega_2,
 \end{array} \right. \quad (4)$$



Recover the temperature on both sides of the membrane, that is $f(t, l_1, y)$ and $p(t, l_1 + \delta_m, y)$, by external boundary measurements.

Thank you